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Structural studies of lithium zinc borohydride by neutron powder diffraction, Raman and NMR spectroscopy

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ABSTRACT

The crystal structure of LiZn₂(BH₄)₅ is studied in detail using a combination of powder neutron diffraction (PND), Raman spectroscopy, and ¹¹B MAS NMR spectroscopy on LiZn₂(BH₄)₅ and LiZn₂(¹¹BD₄)₅. The aim is to obtain detailed structural knowledge of the first interpenetrated framework hydride compound, LiZn₂(BD₄)₅ which consists of doubly interpenetrated 3D frameworks built from dinuclear complex ions $[Zn_2(BD_4)_5]^-$ and lithium ions. The positions of the deuterium atoms are determined using Rietveld refinement of the PND data and the orientation of one of the four independent BD₄⁻ groups is revised. The current data reveal that the structure of $[Zn_2(BD_4)_5]^-$ is more regular than previously reported, as are also the coordinations around the Zn and Li atoms. Both Zn and Li atoms are found to coordinate to the BD₄⁻ units via the tetrahedral edges. Some distortion of the angles within the BD₄ units is observed, relative to the expected angles of 109.4 for the ideal tetrahedral coordination. Raman spectroscopy confirms bending and stretching modes from the expected terminal and bridging bidentate BH_4^- and BD_4^- units. The ¹¹B MAS NMR spectrum of the satellite transitions resolves two distinct manifolds of spinning sidebands, which allows estimation of the ¹¹B quadrupole coupling parameters and isotropic chemical shifts for the four distinct ¹¹B sites of $[Zn_2(BD_4)_5]^-$. These parameters agree favorably well with the isotropic triple-quantum shifts, observed in an ¹¹B MOMAS NMR experiment. The present results may aid in the development and understanding of new borohydride materials for applications as fast ion conductors or as materials for hydrogen storage.

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1. Introduction

Borohydride materials have received much attention recently due to their possible applications as fast ion conductors or as materials for hydrogen storage owing to their high gravimetric hydrogen densities [1–3]. Recently, lithium zinc borohydride has been utilized as a safe store of the poisonous gas diborane [4]. The borohydride chemistry is rich both regarding the reactivity and regarding the chemical bonding, which is mostly ionic for alkali metal borohydrides but more covalent for compounds based on *d*-block metals. However, the thermodynamic and kinetic properties of the metal borohydrides for hydrogen release and uptake are typically less attractive.

Recently, a wide range of novel metal borohydrides have been characterized structurally, physically and chemically and an empirical correlation between the decomposition temperature or the enthalpy of formation and the Pauling electronegativity of the cation was revealed [5]. These studies suggest that the introduction of more covalent bonding in borohydrides facilitates release of hydrogen at lower temperatures and also the formation of more open structures, e.g. Mg(BH₄)₂ [6], Mn(BH₄)₂ [7] and Y(BH₄)₃ [8–10] This relation also holds for a range of new bimetal borohydride, e.g. LiK(BH₄)₂ [11], MSc(BH₄)₄ (M=Li, Na and K) [12–14], NaZn(BH₄)₃, MZn₂(BH₄)₅ (M=Li or Na) [15] as well as the

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first structurally ordered heteroleptic borohydride, KZn(BH₄)Cl₂ [16].

These novel borohydrides exhibit a variety of different types of structures typically forming three dimensional frameworks, but large composite anions are also observed, e.g. $[Sc(BH_4)_4]^-$ and $[Zn(BH_4)Cl_2]^-$ [12-14,16] Moreover, the novel zinc-based borohydrides MZn₂(BH₄)₅ (M = Li or Na) exhibit fascinating novel structural motifs build of doubly interpenetrated 3D frameworks consisting of dinuclear complex anions $[Zn_2(BH_4)_5]^-$ and Li⁺ or Na⁺, acting as counter ions. This structural architecture has previously only been observed for coordination polymers involving organic ligands, also known as metal organic frameworks (MOFs) [17].

In this study, we focus on the detailed structure of the dinuclear complex ion $[Zn_2(BH_4)_5]^-$ in $LiZn_2(BH_4)_5$ combining information from three complementary techniques: powder neutron *diffraction* and Raman and NMR *spectroscopy*. The goal of the present study is to gain insight into the coordination of hydrogen atoms in the structure of the $LiZn_2(BH_4)_5$ and to obtain a more complete description of the metal-hydrogen bonds present in this material. This information may hold the key to obtain a more detailed understanding of the first example of a hydride interpenetrated framework.

2. Experimental

2.1. Sample preparation

All sample handling and preparations were performed under inert argon atmosphere in a MBraun Unilab glove box fitted with a recirculation gas purification system and gas/humidity sensors. Oxygen and water levels were kept well below 1 ppm during all operations.

Sample A (prepared at IFE): LiZn₂(¹¹BD₄)₅ was prepared from a mixture of double isotopically substituted Li¹¹BD₄ (>98%, Katchem) and ZnCl₂ (99.9% Aldrich) in a ratio of 2.5:1 milled in a Spex 6750 freezer-mill at liquid nitrogen temperature for 5 h. The mill was equipped with a stainless steel vial and impactor, and milling was performed with an agitation frequency of about 30 Hz for periods of 15 min each, followed by a 3 min cool-down time.

Samples B: LiZn₂(BH₄)₅ and C: double isotopically substituted LiZn₂(¹¹BD₄)₅ (prepared at Aarhus University) were prepared from mixtures of LiBH₄ (\geq 90%, Aldrich) or Li¹¹BD₄(>98%, Katchem), respectively and ZnCl₂ (\geq 98%, Aldrich) in molar ratio 2:1. Both samples were ball milled in a Fritch Pulverisette No. 4 using the same procedure, i.e. high energy ball milling under inert conditions (argon atmosphere) comprised of 60 times 2 min of milling each intervened by 2 min breaks to avoid heating of the sample and a sample to ball ratio of approximately 1:35. All samples contained LiZn₂(BH₄)₅ or LiZn₂(BD₄)₅, LiCl and small amounts of ZnCl₂ after ball milling.

2.2. Powder neutron diffraction (PND)

PND data were collected for sample **A** with the PUS instrument at the JEEP II reactor at Kjeller, Norway [18]. Neutrons with the wavelength $\lambda = 1.5549$ Å were obtained from a Ge(511) focusing monochromator. The detector unit consists of two banks of seven position-sensitive ³He-detectors, each covering 20° in 2 θ (binned in steps of 0.05°). Data were collected in the 2 θ range 10–130° at T=22°C. The sample was placed in a rotating cylindrical vanadium sample holder (diameter 6 mm).

2.3. Structure refinement

Rietveld refinement was performed using the GSAS program [19]. Pseudo-Voigt functions were applied (constant wavelength profile function type 2), for which the Caglioti (U, V, W) terms, Scherer broadening as well as peak symmetry were refined. The background was modeled by fitting a 30 terms shifted Chebyschev polynomial to manually chosen points. Three different phases were refined for sample A: $LiZn_2(^{11}BD_4)_5$, LiCl (ICSD data file No. 26909) and $ZnCl_2$ (ICSD data file No. 26154). The latter two compounds showed no signs of significant differences from the structures reported in the ICSD files. The Rietveld refinement of the structure concluded that the sample contained about 43 wt% LiZn2(11BD4)5, 41 wt% LiCl, and 16 wt% ZnCl₂. The atomic positions for the heavier atoms: Li, B and Zn obtained from a previous SR-PXD experiment on LiZn₂(BH₄)₅ were taken as a starting point for the Rietveld refinement of the neutron diffraction data [15]. Soft constraints were used on the geometry of the BD₄-tetrahedra, with B–D distances initial set to 1.22 Å, and allowed to vary within 0.03 Å limits. The structure was allowed to relax within 0.2 σ around these starting values. Common isotropic displacement factors were refined for each sort of atoms.

The final refinement of the PND data (see supplementary Fig. S1) gave the following results for $LiZn_2(^{11}BD_4)_5$: space group *Cmca*; lattice parameters:

Table 1

Structural data for the orthorhombic mixed-metal borohydride $LiZn_2(^{11}BD_4)_5$ derived from PND data (this study) and for $LiZn_2(BH_4)_5$ from SR-PXD data.[24].

	From PND	From PXD	
Compound	$LiZn_2(^{11}BD_4)_5$	LiZn ₂ (BH ₄) ₅	
Space group	Cmca	Стса	
a (Å)	8.5936(18)	8.6244 (3)	
b (Å)	17.857 (6)	17.8970 (8)	
c (Å)	15.346 (5)	15.4114 (8)	
Ζ	8	8	
$V(Å^3)$	2354.9(11)	2378.76(19)	
M(g/mol)	233.00	211.93	
<i>T</i> (K)	295	295	

a=8.6031(13), *b*=17.8876(4) and *c*=15.3598(3)Å. The agreement factors are R_{wp} =2.78%, R_p =2.22%, and χ^2 =1.56 (not corrected for background).

2.4. Raman spectroscopy

Raman spectroscopic measurements were performed on the samples B, $LiZn_2(BH_4)_5$ and C, $LiZn_2(BD_4)_5$ using a Renishaw InVia reflex spectrometer with a confocal microscope with a ×20 objective. The excitation laser had a power of \sim 2 mW and a wavelength of λ = 488 nm. An Instec HCS621 V cell was used for inert sample transfer and measurement in a flow of 100 mL/min argon at 1 bar.

2.5. Solid state MAS NMR spectroscopy

Solid-state ¹¹B MAS and MQMAS NMR spectra of LiZn₂(BH₄)₅ and LiZn₂(BD₄)₅ (samples B and C) were obtained on a Varian INOVA-300 (7.05 T) spectrometer using a home-built CP/MAS NMR probe for 5 mm o.d. rotors. The experiments used an ¹¹B rf field strength of $\gamma B_1/2\pi = 75$ kHz, a spinning speed of $\nu_R = 10.0$ kHz, and a 4-s relaxation delay. The MQMAS NMR spectrum was obtained with the three-pulse, z-filter sequence [20], employing $\gamma B_1/2\pi = 75$ kHz for the triple-quantum (3Q) excitation and conversion pulses and $\gamma B_1/2\pi = 25$ kHz for the soft z-pulse, 128 increments in the t_1 dimension, and a 72-step phase cycle. The NMR experiments were performed at ambient temperature using air-tight, end-capped zirconia rotors packed with the samples in a Ar-filled glovebox. The determination of the ¹¹B quadrupole coupling parameters was performed by least-squares fitting of simulated to experimental spinning sideband intensities using the STARS software package [21]. The ¹¹B isotropic chemical shifts are in ppm relative to neat F₃B O(CH₂CH₃)₂.

3. Results and discussion

3.1. Detailed crystal structure of $LiZn_2(^{11}BD_4)_5$

In the present study, PND data is utilized for a detailed structural investigation, which is somewhat complementary to the previous study using synchrotron radiation X-ray powder diffraction (SR-PXD) [15]. Hydrogen (H) is the weakest scatterer of X-rays of all elements, however, deuterium (D) is a strong scatterer of neutrons, i.e. thereby the exact structural positions of deuterium can be determined.

From the Rietveld refinement of the PND data for $LiZn_2(^{11}BD_4)_5$ (sample A) the atomic positions of the heavier Zn atom are found to be in full agreement with the structural information obtained from previous SR-PXD studies [15], in contrast to the positions of Li and B positions which slightly differ from the previous findings. Table 1 presents structural data for $LiZn_2(^{11}BD_4)_5$ compared to previous published SR-PXD results for $LiZn_2(BH_4)_5$ and the revised structure is shown in Fig. 1 (the atomic coordinates and displacement parameters are given in supplementary Table S1). The overall structural topology of $LiZn_2(BD_4)_5$ is preserved consisting of two interpenetrated 3D frameworks,. However a distortion of the previously proposed framework structure is observed from the refinement of the PND data.

The most significant difference between the structures obtained from PND and SR-PXD data, is that the orientation of the borohydride unit containing atom B2 has been revised. In the initial refinement of the published SR-PXD model [15] versus the PND data, very large positional discrepancies for the B2D₄ group were obtained. The refinement resulted in a large difference between the two Zn–D



Fig. 1. The crystal structure of LiZn₂(¹¹BD₄)₅ viewed along (a) [100] and (b) [001] showing the three-dimensional framework. The B–D contacts are shown as solid bonds and BD₄-units are shown as transparent tetrahedra. Zn–B and Li–B contacts are shown as dashed lines. Metal-hydrogen bonds have been omitted for clarity (Zn blue, Li yellow, B brown and D gray).

distances. Furthermore, some of the Li–D distances involving the B2D₄ group were unrealistically short. Hence, the initial refinement suggested that the B2D₄ unit coordinates to Li via the tetrahedral face and to Zn via the tetrahedral vertex. Taking the coordination scheme of the other BD₄ units into account, this was found to be rather unlikely. Therefore, the B2D₄ unit was rotated by 90° around the *c*-axis connecting Zn and B atoms. This resulted in a better fit to the PND data decreasing the agreement from R_{wp} = 3.59%, R_p = 2.77%, and χ^2 = 2.07 to R_{wp} = 2.78%, R_p = 2.22%, and χ^2 = 1.56. Furthermore, this correction gave a much more regular coordination both around Li1 and Zn1. Also, the coordination environments for Zn1 and Zn2 atoms are identical in the revised structure.

Fig. 2 shows the local coordination of the two crystallographically independent Zn-atoms in $[Zn_2(BD_4)_5]^-$ with Zn–D contacts shown as bonds. Zn1 and Zn2 atoms have nearly planar trigonal environment considering the three Zn–B contacts. The plane of the Zn–B-triangle around Zn1 is orthogonal to the similar Zn–B triangle around Zn2.

The coordination mode for the BD₄ units to zinc in $[Zn_2(BD_4)_5]^-$ corresponds to the cation…D₂B coordination scheme, i.e. the BD₄ units share an edge with the zinc atoms. The observed coordination of the BH₄⁻ anion via the edge is typical for metal borohydrides [22]. It is interesting to note that the tentative hydrogen position found



Fig. 2. Local coordination around the two different Zn-atoms in the $[Zn_2(BD_4)_5]^-$ anion. Zn–D and B–D contacts are represented as bonds.

from PXD data for the anions $[Zn(BH_4)_3]_n^{n-}$ in NaZn(BH₄)₃ and $[Zn(BH_4)Cl_2]^-$ in KZn(BH₄)Cl₂ suggests that the coordination mode in these compounds also corresponds to the cation...H₂B scheme [15,16]. Thus, the preferred Zn–BH₄ coordination is likely to be edge sharing (η^2). On the other hand, for $[Sc(BH_4)_4]^-$ in MSc(BH₄)₄ M = Li, Na and K, the PXD data suggest that the coordination mode corresponds to the cation...H₃B scheme; that is, a BH₄⁻ unit has a face of three H atoms oriented towards the Sc (η^3) [12–14]. These results tend to indicate that it is the metal center and not the BH₄⁻ ligand that dictates the observed cation-BH₄ coordination. This suggests that the BH₄⁻ units act as a very flexible ligand, adapting to the coordination centre.

The cation…D₂B coordination scheme in the $[Zn_2(BD_4)_5]^-$ anion results in a six-fold coordination of zinc to deuterium or, equivalently, three-fold coordination of Zn to η^2 -BD₄⁻. The coordination of zinc to deuterium forms trigonal prisms around Zn1 and Zn2 (see Fig. 3(a) and (b)). The BD₄ unit of B3 coordinates to both Zn1 and Zn2 via the two opposite tetrahedral edges. This type of nearly linear coordination of the BH₄ group via the opposite edges is typical for metal borohydrides, such as Be(BH₄)₂ [23], Mg(BH₄)₂ [6,24,25] and Mn(BH₄)₂ [7].

The interatomic distances and angles found in the final refinement (see supplementary Table S2) show a small spread, suggesting that the BD₄ tetrahedra are slightly distorted. Average values of bond lengths and angles are however, in accordance with previous studies of M-BH₄ compounds [22] and close to the expected value of 109.4 for the ideal tetrahedral coordination. While the deviations from the average B-D distance are statistically insignificant, the D-B-D angles in the Zn-D₂B and Li-D₂B fragments are slightly more open than for the non-coordinated D₂B edges. The average D-B-D angles are 112.4 and 115.9 respectively in the Zn-D₂B and Li-D₂B fragments, versus 107.3 for the D-B-D angles of the noncoordinating edges. Similar opening of the coordinated BH2 edges was observed both experimentally by single crystal X-ray diffraction and by DFT calculations in α -Mg(BH₄)₂ (P6₁22) and Al(BH₄)₃ [6,26-28]. While the experimental difference between the coordinated and non-coordinated H-B-H angles in Mg(BH₄)₂ is only 2.0°, the difference is 6.5° in LiZn₂(BD₄)₅. In fact, the D–B–D angles



Fig. 3. (a) Trigonal prism showing the coordination of deuterium to Zn1 viewed along [100]. (b) trigonal prism showing the coordination of deuterium to Zn2 viewed along [0-1-1] and (c) dodecahedron showing the coordination of deuterium to Li1 viewed along [010].

found for LiZn₂(BD₄)₅ by PND correspond very well to the 115(1)° and 107(1)° values obtained by DFT calculations for α -Mg(BH₄)₂ [29]. This reflects that the results obtained by PXD, are affected by the displacement of the electron cloud around the hydrogen atoms towards the metal, which results in observation of smaller opening angle. The positions of the nuclei are observed by PND, i.e. the actual opening angles are determined experimentally, and seem to confirm the results obtained by DFT calculations [29].

The Zn–D bonds from Zn1 are all similar, within the range from 1.799(16) to 1.835(18) Å. The Zn–D bonds involving Zn2 atom cover a wider range from 1.80(2) to 1.94(3) Å. The B-Zn–B angles are also similar, ranging from 116.5(9)° to 126.2(9)°, with the angles between the terminal BD₄⁻ units (B1, B2 and B4) larger than those between the terminal and bridging BD₄⁻ units. The bonding through Zn1–B3D₄–Zn2 define an almost linear coordination with an angle of 176.5(7).

The Li atom has a distorted saddle-like environment by four BD_4^- units, each from a different $[Zn_2(BD_4)_5]^-$ complex anion, hence Li connects the anionic fragments into a 3D framework. As for Zn, Li coordinates to the BD_4 tetrahedra in a bidentate fashion via the tetrahedral edges, i.e. the BD_4^- units act as η^2 -ligands. Thus, Li coordinates to eight D-atoms forming distorted dodecahedra. This coordination is similar to that of Li in three of the four known LiBH₄ phases [30,31]. The Li–D bond lengths cover a very wide range from 1.88(7) to 2.639(11) Å.

As mentioned above, $LiZn_2(BD_4)_5$ consist of two interpenetrated 3D frameworks with no bonds between them. This is further underlined by looking separately at one of the two frameworks, containing 1D voids with coordinates *x*, 0.6014, 0.4090. The void radius, measured as the distance from the center to the nearest atom (deuterium), is 3.7 Å. It is similar in size to the empty voids found in α -Mg(BH₄)₂ [6] and the high-temperature phase of Y(BH₄)₃ [8,9]. In the structure of LiZn₂(BD₄)₅, these voids accommodate the second identical 3D framework.

3.2. Raman spectroscopy

Raman spectra of $LiZn_2(BH_4)_5$ (sample B) and $LiZn(BD_4)_5$ (sample C) are shown in Fig. 4. Table 2 lists the observed Raman



Fig. 4. Raman spectra of $LiZn_2(BH_4)_5$ (sample B) and $LiZn_2(^{11}BD_4)_5$ (sample C) measured at *RT* and with subtracted background.

shifts and the assignment for both LiZn₂(BH₄)₅ and LiZn₂(BD₄)₅. For LiZn₂(BH₄)₅ several bands are observed in the B–H stretching mode region from 2000 to 2600 cm⁻¹. Two maxima are observed at 2100 and 2458 cm⁻¹, corresponding respectively, to bidentate B–H stretching of the terminal B–H₂ and the B–H₂–Zn bridge in the terminal BH₄ units (B1, B2 and B4, Fig. 2) of [Zn₂(BH₄)₅]⁻. Furthermore, additional weaker bands are observed between 2200 and 2400 cm⁻¹, which are assigned to bidentate B–H₂ stretching of the BH₄ unit bridging the two Zn atoms in the complex anion [Zn₂(BH₄)₅]⁻. The observation of terminal B–H₂ stretching underlines that [Zn₂(BH₄)₅]⁻ is an isolated anion, i.e. the coordination towards Li is much weaker. Between 990 and 1450 cm⁻¹ three bands are observed most likely corresponding to B–H bending modes [32,33].

The Raman spectrum of $LiZn_2(BD_4)_5$ (sample A) shows the same vibrations, though shifted to lower Raman frequencies due to the heavier deuterium atoms. Due to the m(H)/m(D) mass ratio of 1/2, the frequencies of the corresponding stretching in the tetrehedral

Table 2

 $Observed \ Raman \ shifts, \ shift \ ratios \ and \ assignments \ to \ specific \ bonds \ for \ LiZn_2(BH_4)_5 \ and \ LiZn_2(BD_4)_5.$

Shift		Ratio	Assignment ^a
LiZn ₂ (BH ₄) ₅ (cm ⁻¹)	$LiZn_2(BD_4)_5 (cm^{-1})$	$\omega_{ m D}/\omega_{ m H}$	
2458	1848	0.752	Bidentate B-H ₂ stretching, terminal BH ₄ units (B1, B2, B4)
2356	1745	0.741	Bidentate B-H ₂ -Zn stretching, bridging BH ₄ unit (B3)
2251	1569	0.697	Bidentate B-H ₂ -Zn stretching, bridging BH ₄ unit (B3)
2100	1487	0.708	Bidentate B-H ₂ -Zn stretching, terminal BH ₄ units (B1, B2, B4)
1431	1030	0.720	B–H bending
1159	825	0.712	B–H bending
995	740	0.743	B–H bending
470	430	_	Zn-B stretching
230	230	-	Zn-Cl vibrations in ZnCl ₂

^a B1, B2, B3 and B4 refer to the distinct BH₄ units shown in Fig. 2.

 BD_4 is affected by the characteristic isotope shift factor of approximately 0.707 for an isotropic species [34]. These factors are found to range from 0.697 to 0.740 (see Table 2), indicating a small degree of anisotropy and that boron and possibly the surrounding Zn atoms take part in these internal vibrations as well.

Furthermore, bands are observed at 470 and 430 cm^{-1} for LiZn(BH₄)₅ and LiZn₂(BD₄)₅, respectively. These are assigned to Zn–B stretching, which corresponds well with the small observed isotope shift of this band, indicating that hydrogen and deuterium atoms take a small part in these vibrations. Similar Raman bands are observed for MSc(BH₄)₄ (M = Li, Na and K), which were assigned to Sc-B stretching modes [12–14]. Observation of such bands is indicative of isolated metal-borohydride complexes, since borohydrides containing no terminal BH₄ units e.g. Mn(BH₄)₂ and Mg(BH₄)₂ exhibit no Raman bands in this region [7,35]. The bands observed at 230 cm⁻¹ for both samples are most likely related to Zn–Cl vibrations of the Zn–Cl₄ tetrahedra in ZnCl₂, which is present in small amounts in both samples.

3.3. Solid-state MAS NMR

The ¹¹B MAS NMR spectrum of the ¹¹B (I=3/2) satellite transitions $(m = \pm 1/2 \leftrightarrow m = \pm 3/2)$ for LiZn₂(BD₄)₅ (Fig. 5a) resolves two manifolds of spinning sidebands (ssbs), whereas three distinct peaks are observed in the isotropic dimension (F1) of the ¹¹B MQMAS NMR spectrum (Fig. 6). A ¹¹B MQMAS NMR spectrum of the non-deuterated sample (LiZn₂(BH₄)₅, sample B, not shown), obtained with ¹H decoupling in both the triple-quantum (3Q) evolution and single-quantum (1Q) detection periods, reveals the same resonances as observed in Fig. 6. However, the resolution is slightly better for the deuterated sample, which reflects that the strong ¹¹B-¹H dipolar couplings are not completely eliminated by MAS and ¹H decoupling under the present experimental conditions. Thus, the subsequent analysis focuses on the acquired spectra for $LiZn_2(BD_4)_5$. Examination of the ssb intensities in the MAS NMR spectrum clearly reveals that the manifold at lowest frequency exhibits the smallest quadrupole coupling and an overall intensity corresponding to roughly 20%. Furthermore, the line shapes of the ssbs in the manifold at high frequency include several shoulders and distinct peaks, indicating that this manifold is composed of resonances from two or more ¹¹B sites. Utilizing these distinct features in a least-squares fit of simulated to experimental ssb intensities for the satellite transitions, assuming that the high-frequency manifold contains resonances from three ¹¹B sites in a 1:1:2 ratio, results in the simulated MAS spectrum shown in Fig. 5b and the ¹¹B quadrupole coupling and isotropic chemical shift parameters listed in Table 3. Obviously, an unambiguous determination of these parameters is only achieved for the lowfrequency site, as a result of the overlap of ssbs from the ¹¹B sites constituting the high-frequency ssb manifold. However, the ¹¹B parameters derived from the satellite transitions reproduce very well the isotropic triple-quantum chemical shifts (δ_{30}^{exp}) observed



Fig. 5. (a) Experimental ¹¹B MAS NMR spectrum of $\text{LiZn}_2(\text{BD}_4)_5$, illustrating the manifold of ssbs from the satellite transitions, recorded at 7.05 T using a spinning speed of $U_R = 10.0$ kHz. (b) Optimized simulation of the ssb manifolds in (a), resulting in the ¹¹B quadrupole coupling parameters and isotropic chemical shifts listed in Table 3. The insets illustrate the splitting of the ssbs into two distinct manifolds.

in the ¹¹B MQMAS NMR spectrum (Fig. 6). These shifts can be calculated according to the relation [36]:

$$\delta_{3Q}^{calc} = \frac{17}{8} \delta_{iso} + \frac{1}{32} \frac{C_Q^2 (1 + \eta_Q^2/3)}{\nu_L^2} \times 10^6 \tag{1}$$

where v_L is the ¹¹B Larmor frequency ($v_L = 96.24$ MHz at 7.05 T). The quadrupole coupling parameters are defined as $C_Q = (eQ/h)V_{zz}$ and $\eta_Q = (V_{yy} - V_{xx})/V_{zz}$, where Q is the nuclear quadrupole moment and V_{ii} the principal elements of the electric-field gradient ten-

Table 3

¹¹B Quadrupole coupling parameters (C_Q , η_Q), isotropic chemical shifts (δ_{iso}), and triple-quantum chemical shifts (δ_{3Q}) determined for LiZn₂(BD₄)₅^a.

$\delta_{ m iso}$ (ppm)	$C_{\rm Q}(\rm kHz)$	η_{Q}	Intensity	Assignment ^b	δ^{calcc}_{3Q} (ppm)	$\delta_{3Q}^{\exp d}$ (ppm)
-42.3	945	0.09	1	B1	-87.0	-87.6
-43.1	748	0.32	1	B2	-89.7	-89.5
-43.3	879	0.06	2	B4	-89.4	-89.5
-46.0	553	0.19	1	B3	-96.7	-89.3

^a The C_Q , η_Q , and δ_{iso} values are obtained from simulations of the ssb manifolds in Fig. 5a, assuming the presence of four ¹¹B sites with relative intensities of 1:1:1:2 (see text for further details).

^b Assignment of the resonances to the four B sites in LiZn₂(BD₄)₅, reported from the PND data (Table S1).

^c Isotropic triple-quantum chemical shifts, calculated from the C_Q , η_Q , and δ_{iso} values using Eq. (1).

^d Isotropic triple-quantum chemical shifts, determined from the ¹¹B MQMAS spectrum in Fig. 6.



Fig. 6. ¹¹B MQMAS NMR spectrum (7.05 T, $v_R = 10.0 \text{ kHz}$) of LiZn₂(BD₄)₅, illustrating the resolution of three distinct ¹¹B sites in the isotropic dimension (F1). The vertical and horizontal spectra correspond to summations over the MQMAS spectrum in the isotropic (F1) and anisotropic (F2) dimensions, respectively. The isotropic 3Q shifts (δ_{30}^{exp}) determined from the vertical summation are listed in Table 3.

sor defined as $|V_{zz}| \ge |V_{yx}| \ge |V_{yy}|$. The experimental and simulated isotropic 3Q shifts are compared in Table 3, which reveal differences below 1 ppm for δ_{3Q} , which are below the error limit for this parameter, considering the uncertainty on the chemical shift reference and the resolution in the F1 dimension of the MQMAS experiment.

The present analysis shows that the ¹¹B MAS and MQMAS NMR spectra are in accord with the presence of four distinct sites in a 1:1:1:2 ratio for the $[Zn_2(BD_4)_5]^-$ anion. The coordination of the individual BD₄⁻ units in this anion, strongly suggests that the low-frequency resonance is assigned to the central BD₄⁻ unit (i.e., B3, Fig. 2), which has a unique coordination involving four B-D-Zn bonds. Furthermore, the relative intensities imply that the B4 site is associated with the δ_{iso} = -43.3 ppm resonance. The geometries of BD4- for the B1 and B2 sites (Table S2) indicate a higher degree of distortion for the BD₄⁻ tetrahedron of the B1 site, as evaluated by the average D-B-D bond-angle deviation from perfect tetrahedral symmetry. Thus, the B1 site is assigned to the δ_{iso} = -42.3 ppm resonance (Table 3) which exhibits the largest ¹¹B quadrupole coupling (tetrahedral distortion) of the two remaining resonances. An improved assignment and further insight into the relation between ¹¹B quadrupole coupling parameters and structural features of the BD_4^- units in the $[Zn_2(BD_4)_5]^-$ anion may potentially be achieved from density functional theory (DFT) calculations of the ¹¹B electric field gradients as demonstrated in a recent ¹¹B MAS NMR and DFT study of inorganic borates [37]. Finally, it is noted that the isotropic chemical shifts determined for LiZn₂(BD₄)₅ (Table 3) is very similar to the values reported for the zinc borhydrides KZn(BH₄)Cl₂ (δ_{iso} = -43.6 ppm) [16] and NaZn(BH₄)₃ (δ = -42.0 ppm and δ = -43.7 ppm in a 9.39 T) [15].

4. Conclusion

The detailed structure of $LiZn_2(BH_4)_5$ and $LiZn_2(BD_4)_5$ is investigated utilizing complementary techniques probing long range order, i.e. of powder neutron diffraction, and the first coordination spheres, i.e. Raman and MAS NMR spectroscopy. These studies

show that the title compound contains a complex $[Zn_2(BH_4)_5]^$ anion, where the BH₄ anion acts as a flexible but directional ligand. The complex anions and lithium cations are organized into a doubly interpenetrated framework. The refinement from PND diffraction data complements the results obtained by the previous SR-PXD studies, however, the coordination around both Li and Zn atoms is found to be more regular than previously depicted. The determination of the precise D-positions, results in revision of the orientation of one of the BD₄⁻ units and shows that the BD₄⁻ units act as η^2 -ligands. Furthermore, it is revealed that the coordinated D₂B edges are slightly more open than the non-coordinated D₂B edges.

The Raman spectra support the PND data showing that there are BH₄ groups strongly coordinated by metal atoms in a bidentate η^2 mode. Furthermore, analysis of ¹¹B MAS and MQMAS NMR spectra of LiZn₂(BD₄)₅ is in accord with the presence of four different boron sites in the [Zn₂(BD₄)₅]⁻ anion. The ¹¹B quadrupole coupling parameters and isotropic chemical shifts, reported for the four boron sites, are tentatively assigned to the different BD₄⁻ units reported from the PND data.

The thorough structural analysis of mixed metal borohydrides reveals the fascinating and complex structural chemistry of these compounds that can be described in terms of coordination chemistry. It gives new ideas for design and synthesis of other novel borohydride materials for applications as fast ion conductors or as materials for hydrogen storage.

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Appendix A. Supplementary data

Supplementary data associated with this article can be found, in the online version, at doi:10.1016/j.jallcom.2010.11.008.

References

- [1] S.-I. Orimo, Y. Nakamori, J.R. Eliseo, A. Züttel, C.M. Jensen, Chem. Rev. 107 (2007) 4111–4132.
- [2] W. Grochala, P.P. Edwards, Chem. Rev. 104 (2004) 1283-1315.
- [3] H. Maekawa, M. Matsuo, H. Takamura, M. Ando, Y. Noda, T. Karahashi, S.-I. Orimo, J. Am. Chem. Soc. 131 (2009) 894–895.
- [4] O. Friedrichs, A. Borgschulte, S. Kato, F. Buchter, R. Gremaud, A. Remhof, A. Züttel, Chem. Eur. J. 15 (2009) 5531–5534.
- [5] Y. Nakamori, K. Miwa, A. Ninomiya, H.-W. Li, N. Ohba, S. Towata, A. Züttel, S.-I. Orimo, Phys. Rev. B 74 (2006) 045126.
- [6] Y. Filinchuk, R. Černý, H. Hagemann, Chem. Mater. 21 (2009) 925–933.
 [7] R. Černý, N. Penin, H. Hagemann, Y. Filinchuk, J. Phys. Chem. C 113 (2009)
- 9003–9007. [8] D.B. Ravnsbæk, Y. Filinchuk, R. Černý, M.B. Ley, D. Haase, H.J. Jakobsen, J. Skib-
- sted, T.R. Jensen, Inorg. Chem. 49 (2010) 3801–3809. [9] C. Frommen, N. Aliouane, S. Deledda, J.E. Fonneløp, H. Grove, K. Lieutenant, I.
- [2] C. Honneller, R. Andranc, S. Deredaa, E. Honneler, R. Grove, R. Ledenan, L. Llamas-Jansa, S. Sartori, M.H. Sørby, B.C. Hauback, J. Alloys Compd. 496 (2010) 710–716.
- [10] T. Jaron, W. Grochala, Dalton Trans. 39 (2010) 160-166.
- [11] E.A. Nickels, M.O. Jones, W.I.F. David, S.R. Johnson, R.L. Lowton, M. Sommariva, P.P. Edwards, Angew. Chem. Int. Ed. 47 (2008) 2817–2819.

- [12] H. Hagemann, M. Longhini, J.W. Kaminski, T.A. Wesolowski, R. Černý, N. Penin, M.H. Sørby, B.C. Hauback, G. Severa, C.M. Jensen, J. Phys. Chem. A 112 (2008) 7551–7555.
- [13] R. Černý, G. Severa, D.B. Ravnsbæk, Y. Filinchuk, V. d'Anna, H. Hagemann, D. Haase, C.M. Jensen, T.R. Jensen, J. Phys. Chem. C 114 (2010) 1357–1364.
- [14] R. Černý, D.B. Ravnsbæk, G. Severa, Y. Filinchuk, V. d'Anna, H. Hagemann, D. Haase, J. Skibsted, C.M. Jensen, T.R. Jensen, J. Phys. Chem. C 114 (2010) 19540-19549.
- [15] D. Ravnsbæk, Y. Filinchuk, Y. Cerenius, H.J. Jakobsen, F. Besenbacher, J. Skibsted, T.R. Jensen, Angew. Chem. Int. Ed. 48 (2009) 6659–6663.
- [16] D.B. Ravnsbæk, L.H. Sørensen, Y. Filinchuk, D. Reed, D. Book, H.J. Jakobsen, F. Besenbacher, J. Skibsted, T.R. Jensen, J. Eur. Inorg. Chem. (2010) 1608–1612.
- [17] T.K. Maji, R. Matsuda, S.A. Kitagawa, Nat. Mater. 6 (2007) 142–148.
- [18] B.C. Hauback, H. Fjellvåg, O. Steinsvoll, K. Johansson, O.T. Buset, J. Jørgensen, J. Neutron Res. 8 (2000) 215–232.
- [19] A.C. Larson, R.B. Von Dreele, General Structure Analysis System (GSAS), Los Alamos National Laboratory Report LAUR (2004) 86–748.
- [20] J.-P. Amoureux, C. Fernandez, S. Steuernagel, S. J. Magn. Reson. Ser. A 123 (1996) 116–118.
- [21] J. Skibsted, N.C. Nielsen, H. Bildsøe, H.J. Jakobsen, J. Magn. Reson. 95 (1991) 88–117.
- [22] Y. Filinchuk, D. Chernyshov, V. Dmitriev, Z. Kristallogr. 223 (2008) 649-659.
- [23] D.S. Marynick, W.N. Lipscomb, Inorg. Chem. 11 (1972) 820-823.
- [24] R. Černý, Y. Filinchuk, H. Hagemann, K. Yvon, Angew. Chem. Int. Ed. 46 (2007) 5765–5767.
- [25] J.H. Her, P.W. Stephens, Y. Gao, G.L. Soloveichik, J. Riissenbeek, M. Andrus, J.C. Zhao, Acta Cryst. B 63 (2007) 561–568.

- [26] B. Dai, D.S. Sholl, J.K. Johnson, J. Phys. Chem. C 112 (2008) 4391-4395.
- [27] S. Aldridge, A.J. Blake, A.J. Downs, R.O. Gould, S. Parsons, C.R. Pulham, J. Chem. Soc. Dalton Trans. (1997) 1007–1012.
- [28] J.P. Soulié, G. Renaudin, R. Černý, K. Yvon, J. Alloys Compd. 346 (2002) 200– 205.
- [29] M.J. van Setten, G.A. de Wijs, M. Fichtner, G. Brocks, Chem. Mater. 20 (2008) 4952–4956.
- [30] Y. Filinchuk, D. Chernyshov, R. Černý, J. Phys. Chem. C 112 (2008) 10579– 10584.
- [31] Y. Filinchuk, D. Chernyshov, A. Nevidomskyy, V. Dmitriev, Angew. Chem. Int. Ed. 47 (2008) 529–532.
- [32] T.J. Marks, J.R. Kolb, Chem. Rev. 77 (1977) 263-293.
- [33] R. Černý, K.C. Kim, N. Penin, V. D'Anna, H. Hagemann, D.S. Sholl, J. Phys. Chem. C 114 (2010) 19127–19133.
- [34] G. Herzberg, Molecular Spectra and Molecular Structure, Infrared and Raman Spectra of Polyatomic Molecules, vol. 2, Krieger Publication Company, Malabar, FL, 1991.
- [35] A. Giannasi, D. Colognesi, L. Ulivi, M. Zoppi, A.J. Ramirez-Cuesta, E.G. Bardají, E. Roehm, M.J. Fichtner, Phys. Chem. A 114 (2010) 2788–2793.
- [36] D. Massiot, B. Touzo, D. Trumeau, J.P. Coutures, J. Virlet, P. Florian, P.J. Grandinetti, Solid State Nucl. Magn. Reson. 6 (1996) 73-83.
- [37] M.R. Hansen, G.K.H. Madsen, H.J. Jakobsen, J. Skibsted, J. Phys. Chem. A 109 (2005) 1989–1997.