

Inorganic Chemistry

Thermal Polymorphism and Decomposition of $Y(BH_4)_3$

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The structure and thermal decomposition of Y(BH₄)₃ is studied by in situ synchrotron radiation powder X-ray diffraction (SR-PXD), ¹¹B MAS NMR spectroscopy, and thermal analysis (thermogravimetric analysis/differential scanning calorimetry). The samples were prepared via a metathesis reaction between LiBH₄ and YCl₃ in different molar ratios mediated by ball milling. A new high temperature polymorph of $Y(BH_4)_3$, denoted β - $Y(BH_4)_3$, is discovered besides the Y(BH₄)₃ polymorph previously reported, denoted α -Y(BH₄)₃. β -Y(BH₄)₃ has a cubic crystal structure and crystallizes with the space group symmetry $Pm\overline{3}m$ and a bisected a-axis, a = 5.4547(8) Å, as compared to α -Y(BH₄)₃, a =10.7445(4) Å ($Pa\overline{3}$). β -Y(BH₄)₃ crystallizes with a regular ReO₃-type structure, hence the Y³⁺ cations form cubes with BH_4^- anions located on the edges. This arrangement is a regular variant of the distorted Y^{3+} cube observed in α -Y(BH₄)₃, which is similar to the high pressure phase of ReO₃. The new phase, β -Y(BH₄)₃ is formed in small amounts during ball milling; however, larger amounts are formed under moderate hydrogen pressure via a phase transition from α - to β -Y(BH₄)₃, at ~180 °C. Upon further heating, β -Y(BH₄)₃ decomposes at ~190 °C to YH₃, which transforms to YH₂ at 270 °C. An unidentified compound is observed in the temperature range 215-280 °C, which may be a new Y-B-H containing decomposition product. The final decomposition product is YB₄. These results show that boron remains in the solid phase when Y(BH₄)₃ decomposes in a hydrogen atmosphere and that Y(BH₄)₃ may store hydrogen reversibly.

Introduction

The transition toward a sustainable and environmentally friendly energy system capable of meeting the increasing energy demands is considered one of the greatest challenges in the 21st century. Hydrogen is suggested as a future carrier of renewable energy, but a safe, compact and efficient hydrogen storage method still remains to be identified.¹⁻³ Borohydride-based materials currently receive increasing interest as potential hydrogen storage systems because of their high theoretical hydrogen densities. Unfortunately, many of the well-known borohydrides exhibits poor thermodynamic and

kinetic properties, which hamper their utilization in technological applications.^{4–6}

Recently, several new mixed-metal borohydrides have been synthesized and characterized, and some of these have improved thermodynamic and kinetic properties, for example, LiSc(BH₄)₄ and NaSc(BH₄)₄ are found to decompose in the temperature ranges 140 to 260 °C and 167 to 267 °C, respectively.^{7,8} These are significantly lower temperatures as compared to the corresponding ones for the alkali borohydrides, LiBH4 and NaBH4, which decompose at approximately 380 and 400 °C, respectively.9

Furthermore, recent studies of metal borohydrides reveal a remarkable diversity in the structural topologies, for example, the compounds $LiZn_2(BH_4)_5$ and $NaZn_2(BH_4)_5$ are built

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from two identical interpenetrated three-dimensional frameworks consisting of isolated complex anions, $[Zn_2(BH_4)_5]^-$, whereas $NaZn(BH_4)_3$ consists only of a single three-dimensional network containing polymeric anions of $[Zn(BH_4)_3]_n^{n-.10}$ Isolated complex anions have also been observed in other structures, such as $[Sc(BH_4)_4]^-$ in $MSc(BH_4)_4$ (M = Li or Na) or [Zn- $(BH_4)Cl_2]^-$ in the heteroleptic borohydride, $KZn(BH_4)Cl_2$.^{7,8,11} Thus, studies of structural, physical, and chemical properties of novel metal borohydrides may hold the key to gain further insight and tailor the stability and property of borohydrides.

Besides the thermodynamic and kinetic properties, it is also of vital importance for possible applications of borohydrides as reversible hydrogen carriers to address the chemical decomposition reactions. For some metal borohydrides a release of gaseous boranes such as diborane has been observed, which hampers reversible hydrogen storage.^{12,13} Therefore, formation of metal borides during decomposition is an important indication of possible reversibility, since boron is kept in the solid state. Furthermore, formation of metal borides is preferred to formation of elemental boron because of the high stability of the latter. It appears that formation of metal borides is facilitated by a hydrogen partial pressure of 1 to 5 bar, which is also the working conditions in most possible applications.¹⁴ Therefore, synthesis, structural characterization, and studies of physical and chemical properties of new metal borohydrides may open new routes for design and preparation of hydrogen storage materials, potentially with improved thermodynamic and kinetic properties for hydrogen release and uptake.

Here we report on the synthesis, characterization, and thermal decomposition of Y(BH₄)₃, studied by in situ synchrotron radiation powder X-ray diffraction (SR-PXD), ¹¹B magic-angle spinning (MAS) NMR, thermogravimetric analysis (TGA), and differential scanning calorimetry (DSC). Furthermore, we present the crystal structure of a new $Y(BH_4)_3$ phase denoted β - $Y(BH_4)_3$.

Experimental Section

Synthesis. Anhydrous yttrium chloride, YCl3 (Sigma-Aldrich, 99.99%) and lithium borohydride, LiBH₄ (Sigma-Aldrich, 98%) were mixed in the molar ratios 1:3 and 1:4, and ball-milled (BM) 120 min under inert conditions (argon atmosphere) in a Fritsch Pulverisette planetary mill using 80 mL tungsten carbide steel containers and approximately an 1:35 mass ratio of sample to balls (o.d. 10 mm, tungsten carbide).

To study the decomposition product, a ball milled sample of LiBH₄ and YCl₃(4:1) was mounted in a corundum crucible placed inside a sealed argon filled quartz capillary, and heated in a furnace from room temperature (RT) to 500 °C with a heating rate of 7 °C/min and then kept at a fixed temperature of 500 °C for 90 min.

The preparation and manipulation of all samples were performed in an argon-filled glovebox with a circulation purifier $(p(O_2, H_2O) < 0.1 \text{ ppm}).$

Laboratory X-ray Powder Diffraction. All samples were initially investigated using laboratory powder X-ray diffraction (PXD) to identify the reaction products and estimate the crystallinity of the samples. PXD measurements were performed in Debye-Scherrer transmission geometry using a Stoe diffractometer equipped with a curved Ge(111) monochromator (Cu K α_1 radiation, $\lambda = 1.54060$ Å) and a curved position sensitive detector. Data were collected at RT between 4 and $127^{\circ} 2\theta$ with counting times of \sim 960 s per step. Air-sensitive samples were mounted in a glovebox in 0.5 mm glass capillaries sealed with glue.

In Situ Time-Resolved SR-PXD. One set of in situ timeresolved SR-PXD data was collected at beamline I911-5 of the synchrotron MAX II, Lund, Sweden in the research laboratory MAX-Lab with a MAR165 CCD detector system and a selected wavelength of $\lambda = 0.9077 \text{ Å}^{.15}$ The X-ray exposure time was 30 s. The sample cell was specially developed for studies of gas/ solid reactions and allows high pressure and temperature to be applied. The powdered sample was mounted in a sapphire (Al₂O₃) single-crystal tube (o.d. 1.09 mm, i.d. 0.79 mm) in an argon-filled glovebox ($p(O_2, H_2O) < 1$ ppm). The sample holder was sealed in a glovebox. The temperature was controlled with a thermocouple placed in the sapphire tube 1 mm from the sample.^{16,17} The sample was heated from RT to 500 °C with a heating rate of 7 °C/min.

Another set of SR-PXD data was collected the Swiss-Norwegian Beamlines (SNBL) at the European Synchrotron Radiation Facility (ESRF) in Grenoble, France. A glass capillary (o.d. 0.5 mm) with the sample was heated from RT to 500 °C at a rate of 4 °C/min, while SR-PXD data were collected. The temperature was controlled with a Cyberstar hot air blower. The data were collected using a MAR345 image plate detector at a sample to detector distance of 275 mm, and a selected X-ray wavelengths of $\lambda = 0.822406$ Å. The capillary was oscillated by 20° during exposure of the sample to the X-ray beam for 20 s.

All obtained raw images were transformed to 2D-powder patterns using the FIT2D program,¹⁸ and calibration measurements of the standard NIST LaB₆ sample, masking diffraction spots from the single-crystal sapphire sample holder. Uncertainties of the integrated intensities were calculated at each 2θ -point by applying Poisson statistics to the intensity data, considering the geometry of the detector.

Structure Solution of a New Y(BH₄)₃ Polymorph. Data collected at I911-5, MAX II on the LiBH₄-YCl₃ (4:1) sample at ~180 °C, containing the high-temperature β -Y(BH₄)₃ phase and no trace of the low-temperature α -polymorph, was used for indexing and structure solution. Diffraction peaks were indexed by the program Dicvol²⁰ in a primitive cubic cell with no systematic absences. Thus, the structure solution was attempted in the space group $Pm\overline{3}m$. The structure was solved in this space group by global optimization in direct space using the program FOX,²¹ varying one yttrium position and the orientation of the rigid tetrahedral BH₄ group (B–H distance fixed at 1.22 Å).

The yttrium atom is found at the special position 0, 0, 0 and boron in the BH₄ group at 1/2, 0, 0, thus adopting the structure of the ambient pressure polymorph of ReO3.²² The orientation

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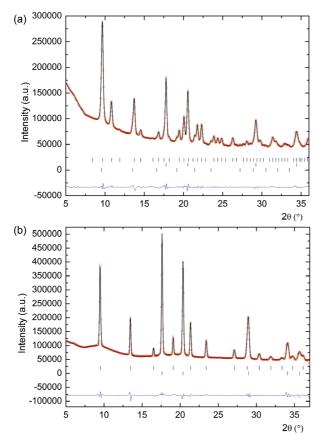


Figure 1. Rietveld plot for the sample of LiBH₄-YCl₃ (4:1) measured at (a) RT and (b) 180 °C. At RT the sample contains a mixture of α -Y(BH₄)₃ (top), β -Y(BH₄)₃ (bottom), and LiCl (middle), whereas the phase transition α - to β -Y(BH₄)₃ has occurred at 180 °C and only β -Y(BH₄)₃ (top) and LiCl (bottom) is present.

of the BH₄ group, however, allows two possibilities: one with hydrogen position at the 24m Wyckoff site x = 0.62910, y =0.12910, z = -0.12910 and the other with hydrogen position at the 24k Wyckoff site x = 0.37095, y = 0, z = 0.18260. The first results in a frustrated net of short (1.86 Å) H···H contacts, which cannot be completely avoided in any, even locally ordered, structure. The second possibility also brings to a net of repulsive (1.45 A) $H \cdots H$ contacts; however, it allows a locally ordered structure to be built with no $H \cdots H$ distances between neighboring BH₄ groups shorter than 2.48 Å. Without distortion or cell doubling, only two subgroups of $Pm\overline{3}m$, namely, Pm3 and P23 allows BH₄ tetrahedra with periodically ordered H-atoms in the 24k Wyckoff site. However, the BH₄ group in the ordered structures would have a square-planar geometry. Thus, it was finally concluded that β -Y(BH₄)₃ has the $Pm\overline{3}m$ structure with BH₄ group disordered over two positions.

The final refinement was performed by the Rietveld method using the program Fullprof.²³ The background was described by linear interpolation between selected points. Only one structural parameter was refined; an overall isotropic atomic displacement, resulting in the excellent refinement factors, $R_{\rm B} =$ 2.6%, $R_{\rm F} = 1.9\%$, $R_{\rm p} = 1.4\%$, $R_{\rm wp} = 2.6\%$ (not corrected for background), and $R_{\rm p} = 8.4\%$, $R_{\rm wp} = 8.3\%$ (corrected for background) and $\chi^2 = 1780$ (this value is high because of the very high counting statistics accumulated by the 2D detector). The Rietveld refinement profile of diffraction from LiCl and β -Y(BH₄)₃ is shown in Figure 1b (data measured at 180 °C) and for a sample containing LiCl, α - and β -Y(BH₄)₃ in Figure 1a (data measured at RT). The sequential refinement in Fullprof was used to refine the cell parameters and weight fractions of the two $Y(BH_4)_3$ polymorphs against the temperature-resolved in situ SR-PXD data.

¹¹**B** MAS NMR. Solid-state ¹¹B MAS NMR spectra were obtained on a Varian INOVA-300 (7.05 T), INOVA-400 (9.39 T), and Direct Drive VNMRS-600 (14.09 T) spectrometers using home-built CP/MAS NMR probes for 5 mm and 4 mm outer diameter rotors. The experiments employed a 0.5 μ s excitation pulse for ¹¹B rf field strengths of $\gamma B_1/2\pi \approx 70$ kHz. The spectra of the Y(BH₄)₃ samples employed ¹H decoupling ($\gamma B_2/2\pi \approx 100$ kHz) during acquisition. The NMR experiments were performed at ambient temperatures using airtight endcapped zirconia rotors, which were packed with the samples in an Ar-filled glovebox. Isotropic chemical shifts are relative to neat F₃B·O(CH₂CH₃)₂, employing a 0.1 M H₃BO₃ aqueous solution ($\delta_{iso} = 19.6$ ppm) as a secondary, external standard sample. Simulations of the MAS NMR spectra were performed using the STARS software package.²⁴

Thermal Analysis. Simultaneous TGA and DSC were performed using a Netzsch STA449C Jupiter instrument heated from RT to 500 °C, $\Delta T/\Delta t = 7$ or 10 °C/min using corundum crucibles with lid as sample holder. The experiments were conducted in helium (4.6) atmosphere.

Results and Discussion

Synthesis and Initial Phase Analysis. Ball milling (BM) cause a metathesis reaction between $LiBH_4$ and YCl_3 , which results in formation of $Y(BH_4)_3$ and LiCl according to eq 1.

$$3LiBH_4 + YCl_3 \rightarrow Y(BH_4)_3 + 3LiCl$$
(1)

The polymorph α -Y(BH₄)₃, previously obtained from diethyl ether solutions of LiBH₄ and YCl₃ at RT and also by BM, is obtained.^{25–27} However, small amounts of a new polymorph of Y(BH₄)₃, denoted β -Y(BH₄)₃, are also formed by BM (see Figure 1a). α -Y(BH₄)₃ can be considered as the stable phase at ambient conditions, whereas β -Y(BH₄)₃ is a high temperature polymorph that can be quenched and is in fact metastable at room temperature. This is similar to the high temperature phases of Mg(BH₄)₂ (space group *Fddd*) and Ca(BH₄)₂ (space group *P*4), which can also be quenched to room temperature.^{28,29}

The fact that the simple metathesis reaction occurs in the mechano-chemical synthesis of $Y(BH_4)_3$ is in contrast to the synthesis of LiSc(BH₄)₄ and NaSc(BH₄)₄, where the mixed-metal borohydrides are formed.^{7,8} It is noteworthy, that scandium and yttrium chlorides react differently even though they are both rare earth ions and have the same charge. However, the difference in ionic radii between that of Sc³⁺ (0.75 Å) and that of Y³⁺ (0.90 Å) might explain this difference in the reaction pathway, that is, the radius of Sc³⁺ is similar to that of Li⁺ (0.76 Å),

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Table 1. Structural Data, Hydrogen Content, Volume Per Formula Unit, and Density for α - and β -Y(BH₄)₃

phase	space group	<i>a</i> [Å]	Ζ	V/Z [Å ³]	$ ho [{ m g}{ m cm}^{-3}]$	$\rho_{v}(H_{2})[kgH_{2}m^{-3})]$	$ ho_{\mathrm{m}}\left(\mathrm{H}_{2} ight)\left[\mathrm{wt}\ \% ight]$
α -Y(BH ₄) ₃	$Pa\overline{\overline{3}}$	10.7445(4)	8	155.2	1.43	129.7	9.07
β -Y(BH ₄) ₃	$Pm\overline{3}m$	5.4547(8)	1	161.5	1.36	123.4	9.07

whereas that of Y^{3+} is significantly larger;³⁰ hence, the coordination of the BH₄⁻ ligand is likely to change.

For the ball milled sample $LiBH_4$ -YCl₃ (4:1) used in this study, no diffraction from the starting compounds is observed by PXD, and from Rietveld refinement the sample composition is found to be 40 wt % α -Y(BH₄)₃, 10 wt % β -Y(BH₄)₃, and 50 wt % LiCl. However, for the sample with the 3:1 composition, both reactants, LiBH₄ and YCl₃, are observed in the milled mixtures by PXD, corresponding to 15 and 30 wt %, respectively, indicating that the metathesis reaction for this composition was not completed under the selected conditions of the BM, which yielded 25 wt % α -Y(BH₄)₃, 4 wt % β -Y(BH₄)₃, and 26 wt % LiCl. For the sample with the 3:1 composition this result was found to be highly reproducible, whereas the amount of reactants found in the ball milled sample LiBH₄-YCl₃ (4:1) varied slightly among the samples.

Comparison of α -Y(BH₄)₃ and β -Y(BH₄)₃ Crystal Structures. The crystal structure of the low-temperature α -Y(BH₄)₃ phase is cubic; space group $Pa\overline{3}$ and cell parameter a = 10.7445(4) Å at RT (see Table 1), in agreement with the previous report.²⁵ The arrangement of the Y^{3+} and BH_4^{-1} ions is that of a distorted ReO₃-type structure, where the BH4⁻ anions are located on the edges of a distorted Y^{3+} cube (see Figure 2a). The mode of BH_4^- coordination by Y^{3+} cations can be determined from the synchrotron X-ray diffraction data only tentatively; however, it is consistent with the model obtained from the neutron diffraction data.²⁵ The coordination mode for the BH₄⁻ anion to yttrium corresponds to the cation \cdots H₂B coordination scheme, that is, the BH₄⁻ anion has an edge of two H-atoms oriented toward the cation. It results in a 12-fold coordination of yttrium to hydrogen or, equivalently, a 6-fold coordination of Y to η^2 -BH₄⁻. The nearly linear coordination of the BH₄ group via the opposite edges is typical in borohydrides, such as $Be(BH_4)_2$,³¹ Mg(BH_4)₂,³² and Mn(BH_4)₂.³³

The new high-temperature β -Y(BH₄)₃ phase also exhibits a cubic crystal structure, with an *a*-axis approximately half as large and in the space group Pm3m (see Table 1). In fact, the structures of α -Y(BH₄)₃ and β -Y(BH₄)₃ are closely related. The structure of β -Y(BH₄)₃ has a regular ReO_3 -type;²² hence, the Y^{3+} cations form a cube with BH4⁻ anions located on the edges (see Figure 2b). This arrangement is a regular variant of the distorted Y^{3+} cube observed in α -Y(BH₄)₃, which is comparable to the high pressure phase of ReO_3 ,³⁴ which shows also a $2 \times 2 \times 2$ cubic superstructure (space group $Im\overline{3}$), however, with a regular Re^{6+} cube and tilted $\operatorname{ReO}_{6}^{1}$

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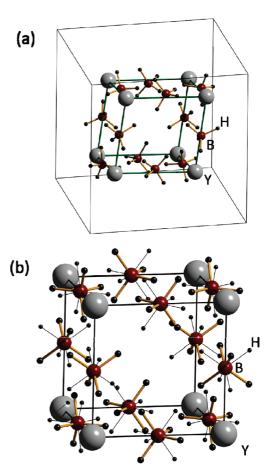


Figure 2. Crystal structure of (a) the low-temperature α -Y(BH₄)₃ phase and (b) the high-temperature β -Y(BH₄)₃ phase. Two orientations of the disordered BH₄ group in the latter structure are depicted by bold and thin lines: a possible local ordering of the BH4 groups, avoiding H···H distances shorter than 2.48 Å, is shown.

octahedra. This means that the octahedral coordination of the Y^{3+} cations to six BH_4^- anions is ideal in β -Y(BH₄)₃ with 6 identical Y–B distances of 2.7274(4) Å at RT (versus 2.714(11) and 2.732(11) Å in the α -phase), 3 identical B-Y-B angles of 180°, and 12 B-Y-B angles of 90° (versus 168.82(3)° and 78.32(3)-104.31(4)°, respectively, in the α -phase). The connection between the YB₆ octahedra via the corner sharing also gives rise to ideal Y-B-Y angles of 180° in the β -Y(BH₄)₃ structure $(168.82(3)^{\circ} \text{ in the } \alpha\text{-phase}).$

It is noteworthy, that yttrium atoms in the α - and β -Y(BH₄)₃ structures have the same coordination number (CN = 12) with respect to the hydrogen atoms as scandium in the recently discovered $MSc(BH_4)_4$ (M = Lior Na) series of borohydrides,^{7,8} despite the difference in ionic radii of Y^{3+} and Sc^{3+} . However, the number of the BH4 units coordinated by the metal atoms and the $M \cdots BH_4$ coordination modes are very different: in MSc- $(BH_4)_4$ (M = Li and Na) four BH_4^- anions are coordinated by Sc via the faces, that is, cation \cdots H₃B coordination scheme or, equivalently, 4-fold Sc to

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 η^3 -BH₄⁻ coordination. Hydrogen forms a distorted cuboctahedron around Sc³⁺ cations in the *M*Sc(BH₄)₄ (*M* = Li or Na) stuctures,^{7,8} while in both the α - and β -Y(BH₄)₃ structures six BH₄⁻ anions are coordinated via the edges (cation···H₂B coordination scheme) forming an isocahedron around the Y³⁺ cations. For the α -Y(BH₄)₃ the isocahedron is somewhat distorted while it is regular for β -Y(BH₄)₃ (Y–H distance is 2.2553(3) Å at RT).

As was pointed out in the Experimental Section, β -Y(BH₄)₃ has *Pm*3*m* symmetry with BH₄ groups disordered over the two positions. The short-range order of the BH₄ group is presumably reducing the number of repulsive $H \cdot \cdot \cdot H$ contacts locally to zero, so that they exceed 2.48 Å. This situation is somewhat similar to the disordered cubic MBH_4 (M = Na-Cs), where the number of the repulsive $H \cdots H$ contacts is minimized via a local ordering of the BH₄ groups.³⁵ The periodic order of the tetrahedral BH_4 groups without short $H \cdots H$ contacts can be achieved in β -Y(BH₄)₃ only in the 2 × 2 × 2 supercell. In the α -Y(BH₄)₃ structure, BH₄ ordering also induces some reorientation of the BH₄ groups and distortion of the YB₆ cubes. These distortions allow the lowtemperature α -Y(BH₄)₃ structure to be more densely packed. The unit cell volume of Y(BH₄)₃ increases drastically on the transition from α -Y(BH₄)₃ to β -Y(BH₄)₃; thus, the high-temperature β -Y(BH₄)₃ structure becomes much more open. This is evident from Figure 6, which shows the unit cell volumes as a function of temperature for the two phases, and reveal that β -Y(BH₄)₃ is 4.5 to 4.7% less dense than α -Y(BH₄)₃ at all temperatures.

The structure of β -Y(BH₄)₃ appears to be porous by observations of a large unoccupied void (~39 Å³) at the position 1/2, 1/2, 1/2, with a distance from its center to the nearest atom (hydrogen) of 3.3 Å. It is similar in size to the void found in the low-temperature phase of Mg(BH₄)₂.³² However, in the small unit cell of β -Y(BH₄)₃ it occupies almost a quarter of the volume. The structure of α -Y(BH₄)₃ is more dense, and the distance from the center of the empty void to the nearest atom (hydrogen) is approximately 2.6 Å; hence, the related volume of the void is much smaller (~11 Å³). This allows us to conclude that the high-temperature β -Y(BH₄)₃ structure is highly porous, likely because of the directional metal···BH₄ bonding defining the non-densely packed architecture of this structure.

¹¹B MAS NMR. The ¹¹B MAS NMR spectra of the central and satellite transitions for samples of LiBH₄– YCl₃ (3:1) and (4:1) (Figures 3a and 3b) include resonances from two distinct ¹¹B sites with chemical shifts of –17.3 ppm and –41.2 ppm and somewhat different relative intensities. The resonance at –41.2 ppm originates from LiBH₄³⁶ and thus, the centerband at –17.3 ppm is assigned to the unique ¹¹B site in α -Y(BH₄)₃. Furthermore, the centerband from this phase has a shoulder at lower frequency, as most clearly observed for the LiBH₄–YCl₃ (4:1) sample (Figure 3b), which is ascribed to the small amount of β -Y(BH₄)₃ present in the

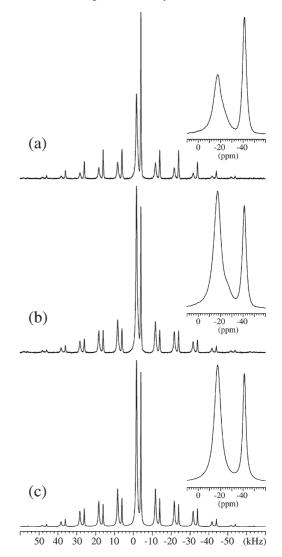


Figure 3. ¹¹B MAS NMR spectra (7.05 T) of samples (a) LiBH₄-YCl₃ (3:1) and (b) LiBH₄-YCl₃ (4:1), recorded with a spinning speed of $\nu_{\rm R}$ = 10.0 kHz, employing high-power ¹H decoupling during acquisition. (c) Simulated spectrum of the resonances from α -Y(BH₄)₃ and LiBH₄ in sample LiBH₄-YCl₃ (4:1), employing the optimized $\delta_{\rm iso}$, $C_{\rm Q}$, and $\eta_{\rm Q}$ parameters listed in the text for these two phases.

sample at ambient temperature as observed from PXD (vide supra). Spectral integration of the resonances from the central and satellite transitions gives the molar fraction of boron present as LiBH₄ and $Y(BH_4)_3$ in the samples. For the LiBH₄-YCl₃ (3:1) sample the molar ratio LiBH₄/ $Y(BH_4)_3$ of 1.00:0.34 is obtained, corresponding to a weight ratio of 1.00:2.05, which is in a good agreement with the weight fractions obtained by Rietveld refinement. For the $LiBH_4$ -YCl₃ (4:1) sample the molar ratio $LiBH_4$ /Y(BH₄)₃ of 1.00: 0.71 (i.e., 1.00:4.37 by weight) is determined which shows that the reaction is more complete since the sample contains a significantly higher quantity of Y(BH₄)₃ as compared to LiBH₄. PXD detects only small amounts of LiBH₄ in the LiBH₄-YCl₃ (3:1) sample, whereas it was not observed for the $LiBH_4$ -YCl₃ (4:1) sample. This indicates that the excess of LiBH₄ is present mainly as an X-ray amorphous phase, most likely caused by the high-energy BM process. Furthermore, the BM may have produced a LiBH₄-LiCl solid solution,³⁶ in which BH_4^- units are partly substituted by Cl⁻ anions (vide infra).

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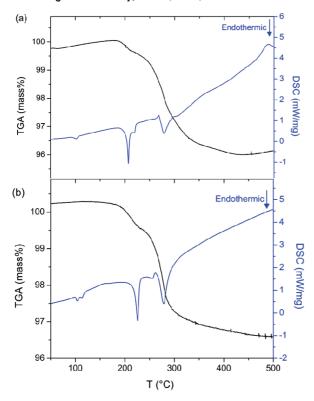


Figure 4. TGA-DSC data for samples containing α -Y(BH₄)₃, LiCl and small amounts of β -Y(BH₄)₃. (a) Ball-milled sample of LiBH₄-YCl₃ (4:1) measured from RT to 500 °C, $\Delta T/\Delta t = 7$ °C/min and (b) ball-milled sample of LiBH₄-YCl₃ (3:1) measured from RT to 500 °C, $\Delta T/\Delta t = 10$ °C/min.

Spinning sidebands from the satellite transitions are observed over a narrow spectral range (<100 kHz), reflecting very small ¹¹B quadrupole couplings for both $Y(BH_4)_3$ and LiBH₄. This is in accordance with previous observations for LiBH₄³⁶ and with the highly symmetric BH_4^- ion in the cubic structure of $Y(BH_4)_3$. The ¹¹B MAS NMR spectra can be satisfactorily simulated using two distinct resonances for α -Y(BH₄)₃ and LiBH₄, as illustrated by the optimized simulation (Figure 3c) of the central and satellite transitions for the $LiBH_4$ -YCl₃ (4:1) sample. The optimization gives the parameters $\delta_{iso} =$ and $\delta_{iso} = -41.2$ ppm, $C_Q = 77.4$ kHz, $\eta_Q = 0.97$ for α -Y(BH₄)₃ and $\delta_{iso} = -41.2$ ppm, $C_Q = 90$ kHz, $\eta_Q = 0.94$ for LiBH₄. The latter values for LiBH₄ are in accordance with the literature.³⁶ A comparison of the experimental and simulated centerbands for $Y(BH_4)_3$ (insets Figure 3) clearly reveals the presence of the low-frequency shoulder at approximately -26 ppm, assigned to the minor quantity of β -Y(BH₄)₃. Thus, the ¹¹B chemical shifts for α - and β -Y(BH₄)₃ are rather similar to those observed for LiSc- $(BH_4)_4$ (-23 ppm³⁷) and NaSc $(BH_4)_4$ (-21.6 ppm⁸).

Thermal Analysis. The ball milled $LiBH_4-YCl_3$ samples are investigated by thermogravimetric analysis/differential scanning calorimetry (TGA/DSC), and the data are displayed in Figure 4. For the sample of molar ratio 4:1 (Figure 4a) the mass loss observed by TGA can be subdivided into two steps: one separate step between 185 and 215 °C, followed by a mass loss covering a wide

temperature range from 215 to 420 °C, the latter likely composed of two or more steps. This indicates a multistep dehydrogenation process. The total observed mass loss in the temperature range 185 to 420 °C is 4.04 wt %, which is less than the calculated hydrogen content (4.28 wt % calculated relative to the sample composition 4:1); hence, no gaseous boranes, $(BH_3)_n$, appear to be lost during the decomposition. This is in agreement with previously reported TGA-DTA and quadrupole mass spectrometry (OMS) data.²⁶

Endothermic peaks are observed by DSC at 101, 207, 219, 260, 278, and 305 °C (peak temperatures). The peak at 101 °C most likely arises from excess LiBH₄ present in the sample, as this temperature corresponds to the phase transition for LiBH₄ from the orthorhombic to the hexagonal high-temperature phase.⁴⁰ The peak at 207 °C corresponds to the first mass loss, whereas the peaks at 260, 278, and 305 °C are observed in the same temperature range as the second mass loss. The latter peak has a very long tail, which may reflect a small mass loss taking place from approximately 330 to 420 °C.

The TGA/DSC data for the sample LiBH₄-YCl₃ in molar ratio 3:1 (Figure 4b) is very similar to the data for the sample with molar ratio 4:1; however, small differences are observed. The mass loss is taking place in two abrupt steps in the temperature ranges from 160 to 230 °C and 230 to 300 °C, followed by a minor slow mass loss occurring from ~310 to 500 °C. The observed mass loss of 4.64 wt % is exactly equal to the calculated hydrogen content (relative to the sample composition 3:1), which suggests that only hydrogen is released from the 3:1 sample. In addition to the peaks observed for the 4:1 sample, a small endothermic peak is observed at 113 °C, whereas the peak at 305 °C is not observed for the 3:1 sample in contrast to the 4:1 sample. The latter suggest that the decomposition pathway is composed of fewer steps for the sample of LiBH₄-YCl₃ in the molar ratio 3:1, which is also in agreement with the separate mass loss steps. These observations for the decomposition of the 4:1 and 3:1 samples will be assigned in the following.

Decomposition Mechanism for Y(BH₄)₃. Figure 5 shows in situ SR-PXD data for the sample consisting of $LiBH_4$ -YCl₃ (4:1). In the temperature range from RT to ~160 °C, diffraction from LiCl, α -Y(BH₄)₃ and small amounts of β -Y(BH₄)₃ are present, and no significant changes in the diffracted intensities are observed. The diffracted intensity from α -Y(BH₄)₃ decreases upon further heating and disappears at ~ 180 °C. In the same temperature interval an increase in the diffracted intensity from β -Y(BH₄)₃ is observed (inset in Figure 5). This suggests that a phase transition from α - to β -Y(BH₄)₃ takes place in the temperature range 160 to 180 °C (at a heating rate of 7 °C/min). β -Y(BH₄)₃ exists in a narrow temperature range from ~ 180 to ~ 190 °C where diffraction is observed to vanish because of decomposition with a mass loss of approximately 0.5 wt % according to the

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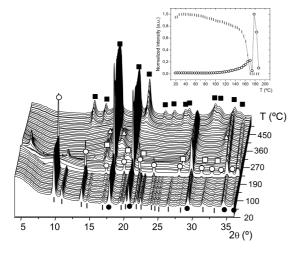


Figure 5. In situ SR-PXD data for LiBH₄-YCl₃ in molar ratio 4:1 ($\lambda = 0.9077$ Å, RT to 500 °C, 7 °C/min). Inset: Normalized integrated intensity of selected Bragg peaks from α - and β -Y(BH₄)₃ as a function of temperature. Symbols: | α -Y(BH₄)₃, \bullet LiCl, $\bigcirc \beta$ -Y(BH₄)₃, \square YH₃, \triangle YH₂, \blacksquare YB₄ (unmarked peaks corresponds to unidentified phases).

TGA analysis. This phase transition is not observed by DSC, although a phase transition has previously been reported by DSC measurements at 201 °C.²⁶ This may indeed correspond to the phase transition from the low-temperature α -Y(BH₄)₃ to the high-temperature polymorph β -Y(BH₄)₃, since LiCl was removed from the sample after BM by use of diethyl ether,²⁶ resulting in a higher quantity of Y(BH₄)₃ implying that the peak from the phase transition is more easily resolved. Furthermore, the chemical separation may slightly alter the thermal properties of Y(BH₄)₃, which may explain the observed temperature discrepancy between the previous DSC results and the PXD data presented in this study.

In the 20–90 °C range, the linear thermal expansion coefficient is slightly lower for the β -phase (1.1 × 10⁻⁵ K⁻¹) than for α -Y(BH₄)₃ (1.4 × 10⁻⁵ K⁻¹). At higher temperatures (100–160 °C) it increases by more than a factor of 2 for both phases, reaching 2.5 × 10⁻⁵ K⁻¹ for the β -phase and 3.1 × 10⁻⁵ K⁻¹ for α -Y(BH₄)₃ (see Figure 6).

YH₃ forms at ~200 °C which suggests that hydrogen is released from the sample. This corresponds well with the temperature for the first mass loss observed by TGA/ DSC. As the temperature reaches ~270 °C diffraction from YH₃ decreases slowly and YH₂ is formed. The two degrees of freedom, temperature and hydrogen content, result in that the peaks of YH₃ and YH₂ exhibit shifts in 2θ positions and are rather broad because of the changes in composition and due to thermal expansion.

At \sim 380 °C, YH₂ decomposes slowly. Both observations of decomposition for YH₃ and YH₂ correspond to the TGA-DSC data in which mass loss and corresponding endothermic peaks are observed at these temperatures, that is, hydrogen is released in these processes. At \sim 400 °C, diffraction from YB₄ slowly emerges. This compound may have formed as an amorphous phase at a lower temperature and crystallizes slowly as the temperature is increased.

Furthermore, in the temperature range from ~ 215 to 280 °C an unidentified compound is observed exhibiting diffraction peaks at 5.25 and 10.60° 2θ (see Figure 5). The formation of this phase is in accord with the endo-

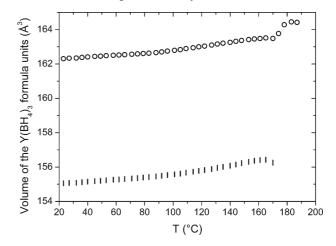


Figure 6. Volume of the Y(BH₄)₃ formula unit (V/Z) for α -Y(BH₄)₃ and β -Y(BH₄)₃ phases as a function of temperature. Symbols: $| \alpha$ -Y(BH₄)₃ and $\bigcirc \beta$ -Y(BH₄)₃.

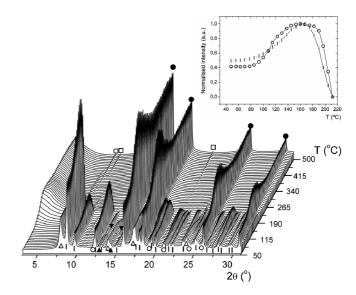


Figure 7. In situ SR-PXD data for LiBH₄-YCl₃ in molar ratio 3:1 ($\lambda = 0.822406$ Å, RT to 500 °C, 4 °C/min). Inset: Normalized integrated intensity of selected Bragg peaks from α - and β -Y(BH₄)₃ as a function of temperature. Symbols: $|\alpha$ -Y(BH₄)₃, \bullet LiCl, $\bigcirc \beta$ -Y(BH₄)₃, \triangle YCl₃, \blacktriangle orthorhombic LiBH₄, \checkmark hexagonal LiBH₄, and \Box Li₃YCl₆ (unmarked peaks correspond to unidentified phases).

thermic peak observed by DSC measurements at 219 °C, while the decomposition may contribute to the peak at 278 °C, also assigned to the transformation of YH₃ to YH₂. This unidentified material may be a new compound in the Y–B–H system similar to other decomposition intermediate phases recently described. Decomposition of LiBH₄ may lead to formation of Li₂B₁₂H₁₂ containing $[B_{12}H_{12}]^{2^-}$ ions as previously observed by X-ray diffraction and solid-state NMR while decomposition of Ca(BH₄)₂ recently revealed a new material CaB₂H₂ containing $[B_2H_2]^{2^-}$ ions.^{37–39}

The in situ SR-PXD data for the sample of LiBH₄– YCl₃ (3:1), displayed in Figure 7, reveals a somewhat different reaction pattern. As previously mentioned, the metathesis reaction was not completed under the selected BM conditions, that is, diffraction from both reactants, LiBH₄ and YCl₃ and products LiCl, α -Y(BH₄)₃, and β -Y(BH₄)₃ are observed. At ~100 °C the phase transformation of the orthorhombic to hexagonal LiBH4 is observed.⁴⁰ As the temperature reaches \sim 90 °C, the amount of LiBH₄ and YCl₃ decreases; meanwhile an increase in the diffracted intensity from LiCl, α -Y(BH₄)₃, and β -Y(BH₄)₃ is observed, that is, the metathesis reaction 1 continues during heating of the sample. The temperature at which this process is initiated may be in accord with the additional DSC peak observed at 113 °C for the 3:1 sample. It is interesting to note that β -Y(BH₄)₃ forms at a higher rate than α -Y(BH₄)₃ during this process (inset Figure 7). Furthermore, β -Y(BH₄)₃ seems to decompose at a temperature ~ 10 °C higher as compared to α -Y(BH₄)₃ form, which is in accordance with the observation for the 4:1 sample.

At ~200 °C diffraction from both α -Y(BH₄)₃ and β -Y(BH₄)₃ vanish, which correspond to the observations for the 4:1 sample (Figure 5) and the TGA/DSC data (Figure 4b). However, the phase transition α - to β -Y- $(BH_4)_3$ at ~180 °C, and the formation of the decomposition products YH₃, YH₂, Y-B-H, and YB₄ are not observed in the SR-PXD data. This may be due to different SR-PXD sample conditions; the 4:1 sample was mounted in a sample holder designed to withstand high pressure, that is, the released hydrogen during the decomposition provide increasing hydrogen pressure, which alters the decomposition pathway drastically. On the other hand, the 3:1 sample was mounted in a glass capillary, which may have leaked the hydrogen released from the sample.

Second, for both samples, 3:1 and 4:1, a solid solution of LiBH₄-LiCl may form during BM where 10-20% of LiCl could have dissolved in LiBH₄ by anion substitution in the LiBH₄ lattice forming Li(BH₄)_{1-x}Cl_x.³⁶ This is known to have a stabilizing effect on LiBH₄, and it may also retard the formation of $Y(BH_4)_3$ during the metathesis reaction 1. Apparently, excess LiBH₄ is needed for a relatively fast formation of Y(BH₄)₃ as observed for the 4:1 composition where all YCl₃ is consumed.

For the composition 3:1, some YCl₃ remains in the sample after BM. Clearly, YCl₃ extracts LiBH₄ from the LiBH₄-LiCl solid solution at elevated temperatures (T > 85 °C) and forms Y(BH₄)₃ according to the reaction Scheme 1. Yttrium borohydride decomposes and forms apparently an amorphous Y-B-H containing phase which may contribute to the significant increase in the background intensity at $2\theta > \sim 7^{\circ}$.

Two unidentified compounds are observed to crystallize after the decomposition of $Y(BH_4)_3$, which are different from those observed in the (4:1) sample. A compound observed in the temperature range ~ 190 to 235 °C may have a composition similar to Y(BH₄)Cl₂ or Y(BH₄)₂Cl previously described.^{41,42} Anion substitution may also have a stabilizing effect in Y(BH₄)₃, resulting in decomposition at higher temperatures. The second phase could be a Y-B-Cl containing decomposition product. At ~375 °C a small amount of Li₃YCl₆ forms, likely in an addition reaction between unreacted YCl₃ and LiCl, as shown in eq 2.

$$YCl_3 + 3LiCl \rightarrow Li_3YCl_6$$
(2)

Additional information about the Y(BH₄)₃ decomposition product is obtained from the ¹¹B MAS NMR

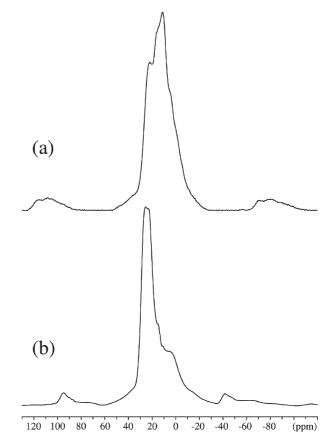


Figure 8. ¹¹B MAS NMR spectra of the LiBH₄-YCl₃ (4:1) sample, heated from RT to 500 °C (7 °C/min.) and stored at this temperature for 90 min, recorded at (a) 9.39 T and (b) 14.09 T using spinning speeds of $\nu_{\rm R}$ = 12.0 kHz and $\nu_{\rm R}$ = 13.0 kHz, respectively.

spectra of the material resulting from heating a bulk sample of $LiBH_4$ -YCl₃ (4:1) in an argon atmosphere to 500 °C (7 °C/min.) and storing it at this temperature for 90 min. (i.e., a heat treatment similar to the in situ SR-PXD experiment, considering the larger amount of sample). The signal intensities, observed in the ¹¹B MAS NMR spectra at two magnetic field strengths (Figure 8), reveal that the decomposition product contains a significant quantity of boron. The absence of narrow spinning sidebands from the satellite transitions indicates that a major part of the boron is present in less-crystalline phases and exhibits quite large ¹¹B quadrupole couplings. The latter is also apparent from the width of the central transitions (Figure 8), which is approximately 60 ppm at 14.09 T. PXD measurement of the same sample shows the presence of YB_4 as the only crystalline boron-containing phase. The ¹¹B quadrupole coupling parameters for the three different boron sites in YB_4 has earlier been reported by static-powder ¹¹B NMR and found to be in accordance with the corresponding ${}^{11}B$ electric-field gradients, obtained by density functional theory calculations.⁴³ These data shows that YB₄ exhibits rather small quadrupole couplings, that is, $C_{\rm O} = 1.03$, 1.40, and 1.03 MHz for B(1), B(2), and B(3), respectively. This will result in rather narrow resonances for the ¹¹B central transitions at 9.39 and 14.09 T and thus, YB_4 may only account for the narrow peaks at 25.8 and 22.7 ppm in the 14.09 T spectrum (Figure 8b) and in the range

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22.3–11.1 ppm at 9.39 T (Figure 8a) and thereby not the dominating part of the central-transition intensity in the ¹¹B MAS NMR spectra. The chemical shifts for the boron sites in the decomposition product is expected to be above ~20 ppm, considering the strong quadrupolar broadening, which imply that borate or borohydride species are not present in the sample. Thus, we ascribe the main part of the intensity in the ¹¹B MAS NMR spectra to less-crystalline boron phase(s) including multiple ¹¹B sites, and potentially containing yttrium. We note that an ¹¹B MQMAS NMR spectrum (14.09 T) of the decomposition product only resolved two resonances, most likely arising from the small quadrupole couplings in the YB₄ phase.

Conclusion

The mechanism for the thermal decomposition of α - and β -Y(BH₄)₃ is studied by in situ synchrotron radiation PXD, ¹¹B MAS NMR spectroscopy, and thermal analysis (TGA/DSC). A new high-temperature polymorph of Y(BH₄)₃, denoted β -Y(BH₄)₃, is discovered, which has lower density and higher crystal symmetry of the cubic crystal structure as compared to the known phase α -Y(BH₄)₃. We have demonstrated that a moderate hydrogen pressure may facilitate stabilization of boron as a metal boride in the solid dehydrogenated phase. These results are very promising for attaining the reversible hydrogen storage in Y(BH₄)₃. Yttrium borohydride has an α - to β -Y(BH₄)₃ phase transition at ~180 °C, β -Y(BH₄)₃ at ~190 °C to YH₃, which transforms to YH₂ at

270 °C. An unidentified phase is observed in the temperature range 215 to 280 °C, which may be a new Y–B–H containing decomposition product. The final decomposition products are YB₄ as well as amorphous boron containing compounds. Thermogravimetric measurements suggest that hydrogen solely is released during the decomposition in the temperature range 185 to 420 °C, which is in strong contrast to LiBH₄ that decomposes at T > 400 °C. The ¹¹B MAS NMR further underlines that boron was retained in the solid phase during decomposition. This study reveals that the stability and decomposition reactions mechanism can be tailored by the metal in the crystal structure and the gas atmosphere.

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Supporting Information Available: X-ray crystallographic files in CIF format. This material is available free of charge via the Internet at http://pubs.acs.org.